Dynamics of branes in doubled space

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based on works with

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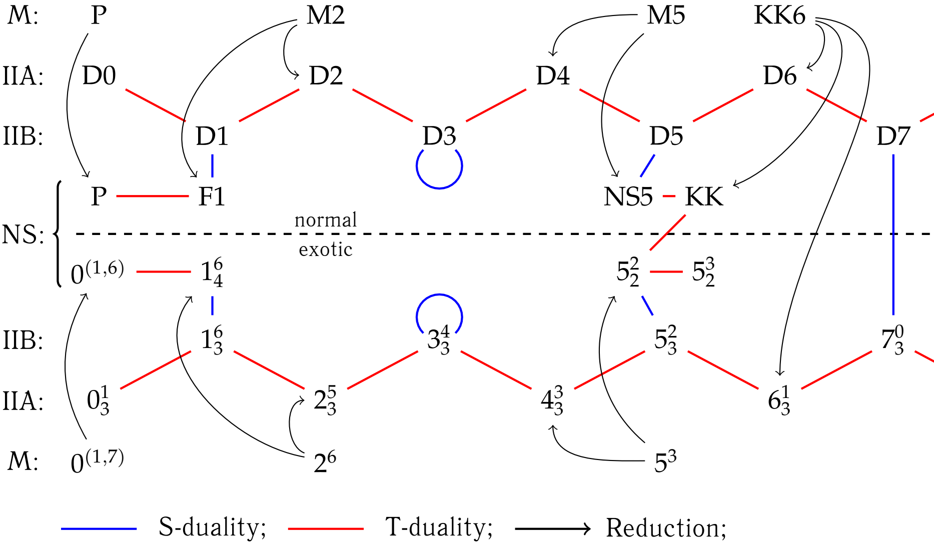
QFTG
Tomsk, 2018

Web of (some) branes



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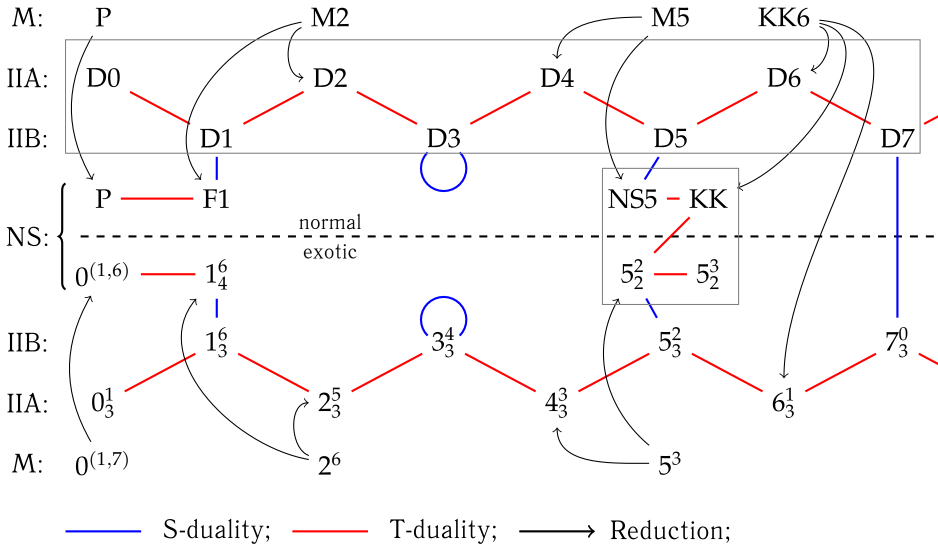

Why bother?

D-branes

- AdS/CFT correspondence: different limits of a single D-brane effective action;
- phenomenological constructions: braneworld models, interesection of branes;
- black hole entropy counting;

NS-branes

- Little String Theories: $\mathcal{N} = (2, 0)$ (for IIA) and $\mathcal{N} = (1, 1)$ (for IIB) 6D theories;
- non-conformal field theory holographies;
- cosmological moduli stabilization via (non-geometric) NS-NS fluxes;

The results

T-duality orbits

$$\mathrm{D0} - \mathrm{D1} - \mathrm{D2} - \mathrm{D3} - \mathrm{D4} - \mathrm{D5} - \mathrm{D6} - \mathrm{D7} - \mathrm{D8} - \mathrm{D9}$$

$$\mathrm{NS5}(5_2^0) - \mathrm{KK5}(5_2^1) - \mathrm{Q}(5_2^2) - \mathrm{R}(5_2^3) - \mathrm{R}'(5_2^4)$$

- Effective actions for these T-duality orbits has been constructed.
- Depending on orientation these project down to actions for normal branes.
- One observes non-geometric effects for D-branes

D-branes: geometry and dynamics

In supergravity D-branes

- are RN black-hole-like solutions of SUGRA equations of motion,
- preserve $\frac{1}{2}$SUSY
- are described by non-zero metric $G_{\mu\nu}$ and gauge field $C_{\mu_1\dots\mu_p}$

As fundamental objects D-branes

- are described by effective action

$$S_{DBI} = \int_\Sigma d^{p+1}\xi \sqrt{\det\left(G_{\mu\nu}\partial_a x^\mu \partial_b x^\nu + \mathcal{F}_{ab}\right)} + \int_\Sigma \mathcal{C}_{p+1} \tag{1}$$

- carry worldvolume fields $x^\mu = x^\mu(\xi)$ and $\mathcal{F}_{ab}$

T-duality

- Mass spectrum of a string on a torus $\mathbb{T}^d$ is invariant under $O(d,d)$ group
- SUGRA solutions transform into solutions
- String does not feel the change in backgrounds
- T-duality is performed along isometries

T-duality orbit of NS branes

	0	1	2	3	4	5	6	7	8	9
NS5 :	$\times$	$\times$	$\times$	$\times$	$\times$	$\times$	$\bullet$	$\bullet$	$\bullet$	$\bullet$
	world-volume						transverse			
KK5 :	$\times$	$\times$	$\times$	$\times$	$\times$	$\times$	$\bullet$	$\bullet$	$\bullet$	$\odot$
	world-volume						transverse			special
Q :	$\times$	$\times$	$\times$	$\times$	$\times$	$\times$	$\bullet$	$\bullet$	$\odot$	$\odot$
R :	$\times$	$\times$	$\times$	$\times$	$\times$	$\times$	$\bullet$	$\odot$	$\odot$	$\odot$
R$'$:	$\times$	$\times$	$\times$	$\times$	$\times$	$\times$	$\odot$	$\odot$	$\odot$	$\odot$

NS5 $\leftrightarrow$ KK5: T_9

KK5 $\leftrightarrow$ Q: T_8

(2)

[deBoer, Shigemori]

T-duality orbit of D-branes

$$
\begin{array}{rcccccccccc}
 & 0 & 1 & 2 & 3 & 4 & 5 & 6 & 7 & 8 & 9 \\
\mathrm{D0}: & \times & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet \\
\mathrm{D1}: & \times & \times & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet \\
\mathrm{D2}: & \times & \times & \times & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet \\
\mathrm{D3}: & \times & \times & \times & \times & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet \\
\mathrm{D4}: & \times & \times & \times & \times & \times & \bullet & \bullet & \bullet & \bullet & \bullet \\
\mathrm{D5}: & \times & \times & \times & \times & \times & \times & \bullet & \bullet & \bullet & \bullet \\
\mathrm{D6}: & \times & \times & \times & \times & \times & \times & \times & \bullet & \bullet & \bullet \\
\mathrm{D7}: & \times & \times & \times & \times & \times & \times & \times & \times & \bullet & \bullet \\
\mathrm{D8}: & \times & \times & \times & \times & \times & \times & \times & \times & \times & \bullet \\
\mathrm{D9}: & \times & \times & \times & \times & \times & \times & \times & \times & \times & \times
\end{array}
\tag{3}
$$

! T-duality changes dimension of a D-brane

Covariant potentials

- D-brane potentials $C_{(p+1)}$ can be combined:

$$|\chi\rangle = \sum_{p=0}^{10} C_{m_1 \dots m_p} \Gamma^{m_1 \cdots m_p} |0\rangle \tag{4}$$

- $O(10,10)$ algebra that includes $GL(10)$ as $T^M = (T^m, T_m)$:

$$\begin{aligned} \{\Gamma_M, \Gamma_N\} &= 2\eta_{MN} \quad \Longrightarrow \quad \{\Gamma_m, \Gamma^n\} = \delta_m{}^n, \\ \text{Clifford vacuum:} & \qquad \Gamma_m |0\rangle = 0 \end{aligned} \tag{5}$$

Covariant potentials

Charge $\langle Q|$ – a smart way to single out $C_{(p)}$ from $|\chi\rangle$:

$$\begin{gathered}\text{for say } \langle Q| = \langle 0| \\ \text{the only non-zero: } \langle Q|\Gamma_{m_1\dots m_{10}}|\chi\rangle = C_{m_1\dots m_{10}} \\ \\ \text{for say } \langle Q| = \langle 0|\Gamma_{\hat{m}} \\ \text{the only non-zero: } \langle Q|\Gamma_{m_1\dots m_9}{}^{\hat{m}}|\chi\rangle = C_{m_1\dots m_9}.\end{gathered} \tag{6}$$

Invariant interaction

$$S_{wz} = \int d^{10}\xi \; \varepsilon^{a_1\dots a_{10}} \langle Q|\Gamma_{M_1\dots M_{10}}|\chi\rangle \partial_{a_1} X^{M_1} \cdots \partial_{a_{10}} X^{M_{10}}\,. \tag{7}$$

Upon choice of $\langle Q|$ this reproduces smth like

$$S_{wz}{}^{D} = \int C_{m_1\dots m_p} dX^{m_1} \wedge \cdots \wedge dX^{m_p} \tag{8}$$

Doubled geometry

- Doubled coordinates $X^M = (x^m, \tilde{x}_m)$
- section constraint for consistency of the theory, kills half of the coordinates
- T-duality: $T_x : x \longleftrightarrow \tilde{x}$
- Generalized metric is a T-duality covariant object

$$\mathcal{H}_{MN} = \begin{bmatrix} g - Bg^{-1}B & Bg^{-1} \\ g^{-1}B & g^{-1} \end{bmatrix}, \quad \text{in analogy with} \quad F_{\mu\nu} = \begin{bmatrix} 0 & \vec{E} \\ -\vec{E} & *_3\vec{B} \end{bmatrix}$$

$$\mathcal{H}_{MN} \in \frac{O(10,10)}{O(1,9)\times O(1,9)} \tag{9}$$

- There exists an action for $\mathcal{H}_{MN}$

[Berman, Cederwall, Coimbra, Godazgar[2], Grana, Hohm, Hull, EtM, Nicolai, Perry, Samtleben, Thompson, Waldram, Zwiebach ...]

Dynamics

Dynamics is ruled by DBI action

$$S_p = \int d^{p+1}\xi e^{-\phi}\sqrt{-\det\left(G_{mn}\partial_a x^m \partial_b x^n + \dots\right)}, \tag{10}$$

$x^m = x^m(\xi)$ scalar fields

A T-invariant version then would be

$$S_D = \int d^{10}\xi e^{-d}\sqrt{-\det\left(\mathcal{H}_{MN}\partial_a X^M \partial_b X^N + \dots\right)}, \tag{11}$$

Dynamics

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A T-invariant version then would be

$$S_D = \int d^{10}\xi e^{-d} \det|h_{\alpha\beta}|^{\frac{1}{4}} \sqrt{-\det\left(\mathcal{H}_{MN}\hat{\partial}_a X^M \hat{\partial}_b X^N + \dots\right)}, \tag{11}$$

where one needs the projected derivatives

$$\begin{aligned} \hat{\partial}_a X^M &= \partial_a X^M - (h^{-1})^{\alpha\beta} k_\alpha{}^M k_\beta{}^N \mathcal{H}_{NK} \partial_a X^K, \\ h_{\alpha\beta} &= k_\alpha{}^M k_\beta{}^N \mathcal{H}_{MN}. \end{aligned} \tag{12}$$

Embedding of branes

	0	1	2	3	4	5	6	7	8	9	$\tilde{1}$	$\tilde{2}$	$\tilde{3}$	$\tilde{4}$	$\tilde{5}$	$\tilde{6}$	$\tilde{7}$	$\tilde{8}$	$\tilde{9}$
D0	k	•	•	•	•	•	•	•	•	•	k	k	k	k	k	k	k	k	k
D1	k	k	•	•	•	•	•	•	•	•	•	k	k	k	k	k	k	k	k
D2	k	k	k	•	•	•	•	•	•	•	•	•	k	k	k	k	k	k	k
D3	k	k	k	k	•	•	•	•	•	•	•	•	•	k	k	k	k	k	k
D4	k	k	k	k	k	•	•	•	•	•	•	•	•	•	k	k	k	k	k
D5	k	k	k	k	k	k	•	•	•	•	•	•	•	•	•	k	k	k	k
D6	k	k	k	k	k	k	k	•	•	•	•	•	•	•	•	•	k	k	k
D7	k	k	k	k	k	k	k	k	•	•	•	•	•	•	•	•	•	k	k
D8	k	k	k	k	k	k	k	k	k	•	•	•	•	•	•	•	•	•	k
D9	k	k	k	k	k	k	k	k	k	k	•	•	•	•	•	•	•	•	•

! Depending on the choice of **k**'s one gets different D-branes

!? D-branes can localize in dual space

Embedding of branes

NS five branes localize in dual space as well:

	world-volume						transverse directions							
	x^0	x^1	x^2	x^3	x^4	x^5	y^1	y^2	y^3	y^4	$\tilde{y}_1$	$\tilde{y}_2$	$\tilde{y}_3$	$\tilde{y}_4$
NS5	×	×	×	×	×	×	•	•	•	•	k	k	k	k
KK5	×	×	×	×	×	×	•	•	•	k	k	k	k	•
Q	×	×	×	×	×	×	•	•	k	k	k	k	•	•
R	×	×	×	×	×	×	•	k	k	k	k	•	•	•
R′	×	×	×	×	×	×	k	k	k	k	•	•	•	•

(13)

• — localization direction, k — Killing direction

Non-perturbative instanton corrections on world-sheet of string localize KK-monopole in dual space. The same is true for 5^2_2-branes. [Jensen, Tong, Harvey, Kimura]

The message

- One is able to construct a single action for several branes, related by T-duality
- For D-branes this suggests localization in dual space

Discussion

- Prove microscopically, that D-branes localize in dual space, calculate instanton corrections
- Field theories on worldvolume (especially for D-branes)
- Generalize stuff for exceptional field theories and U-dualities

What's the use of all that?

- Tadpole cancellation conditions for flux compactifications (Bianchi identities), support for internal space
- String behavior on such backgrounds: non-commutativity and non-associativity
- Little string theories from NS five-branes
- New stuff for AdS / CFT correspondence?

Thank you!